Subsolidus phase relations of the Y₂O₃–SrO–CuO system

WU FEI, XIE SISHEN, CHEN ZHAN, LING JINGKUI Institute of Physics, Chinese Academy of Science, Beijing 100080, People's Republic of China

The 960 °C subsolidus phase relations of Y_2O_3 -SrO-CuO system have been determined by X-ray diffraction. The section at room temperature in the Y_2O_3 -SrO-CuO ternary system can be divided into six three-phase and two-phase regions. In this system, there are five binary compounds: Y_2SrO_4 , $Y_2Cu_2O_5$, Sr_2CuO_3 , $SrCuO_2$ and $Sr_{14}Cu_{24}O_{41}$ as well as a solid solution $Y_xSr_{14-x}Cu_{24}O_{41}$ ($0 < x \le 4$). The structure of the solid solution is discussed. The variation of the lattice parameters of the solid solution with increasing yttrium content has been measured. Electrical resistivity measurements of the solid solution show that there is no superconductive phase in the Y-Sr-Cu-O system.

1. Introduction

Since the high-temperature superconductivity was discovered in the La-Ba-Cu-O [1] and Y-Ba-Cu-O [2, 3] systems, much substitutive work has been done. It has been verified that vttrium can be replaced by most rare-earth elements (except Ce, Pr, Pm and Tb) and the superconductivity and basic structure of the superconductive phase can be preserved. Because alkaliearth metals barium and strontium have similar chemical properties and ionic radii, it is significant to study the change in superconductivity and structure when strontium completely replaces barium. Mei et al. [4] reported $Y_{0.3}Sr_{0.7}CuO_{4-y}$ with a T_c of about 40 K. Wu et al. [5] have reported that in $YSrCuO_{4-y}$ there are two superconductive transitions at 80 and 40 K, corresponding to YSr₂Cu₃O₇ (123) and Y₂SrCu₄O₈ (214), respectively. Oda et al. [6] reported that the tetragonal YSr₂Cu₃O₇ exhibited a complete transition in resistance at 81 K. Zhang et al. [7] found that a zero resistance was obtained at 90.5 K for $Sr_xY_{1-x}CuO_{3-y}$ (x = 0.375–0.5). In contrast, Veal et al. [8] reported that samples with the starting composition YSr₂Cu₃O₇ were insulating and showed no characteristics of an orthorhombic YBa₂Cu₃O₇ phase. In order to search for new superconductive phases, and to identify whether $YSr_2Cu_3O_7$ exists, we investigated the phase relations and the possible superconductive phases in the Y2O3-SrO-CuO ternary system, more carefully.

2. Experimental procedure

The starting materials were Y_2O_3 (99.99%), BaCO₃ (AR) and CuO(AR). The samples were prepared by solid state reactions. The raw mixture of these materials with different compositions was ground thoroughly, then pressed into pellets, sintered at 930 °C for 24 h in air, then slowly cooled to room temperature. The samples were then reground, repressed into pellets, sintered at 960 °C in air for further 48 h, then

3082

slowly cooled in the furnace to room temperature for phase identification or electrical resistivity measurements.

The phase identifications were made with a Guinier-de Wolf monochromatic focusing camera by using CuK_{α} radiation. Pure silicon powder was added as internal standard for calibration of the lattice parameters. The diffraction data of Sr-Cu oxides and the solid solution $Y_xSr_{14-x}Cu_{24}O_{41}$ (0 < x < 4) were collected using a D/max-RB type X-ray diffractometer with a rotating anode generator (Japanese Rigaku).

Electrical resistivity was determined by the fourprobe technique. The mutual-inductance method was adopted to measure the a.c. magnetic susceptibility.

3. Results and discussion

The phase diagram in the Y_2O_3 -SrO binary system has been reported previously [9]. In this system, below 1500 °C only one binary compound, Y₂SrO₄, exists. The phase diagram of the Y_2O_3 -CuO binary system has also been reported [10]. In this system only one binary compound, Y₂Cu₂O₅, is formed. According to the JCPDS cards, Y2SrO4 belongs to an orthorhombic system with a = 1.00070 nm, b = 1.19141 nm and c = 0.34098 nm. Its space group is Pnam (62). Y₂Cu₂O₅ belongs to a monoclinic system with space group $P_2(3)$. Its lattice parameters are a = 1.360 nm, b = 0.327 nm, c = 1.360 nm and β = 133.15° . In the system SrO-CuO, we obtained the following compounds: Sr₂CuO₃, SrCuO₂ and Sr₁₄Cu₂₄O₄₁. The X-ray analysis shows that Sr₂CuO₃ belongs to the orthorhombic system, the unit cell is a body centred lattice with a = 12.708 nm, b = 0.3913 nm and c = 0.3502 nm. SrCuO₂ belongs to an orthorhombic system with space group Cmcm (63), its lattice parameters are a = 0.3573 nm, b = 1.6333 nm and c = 0.3915 nm, coinciding with those given by Teske et al. [11, 12]. The compound Sr₁₄Cu₂₄O₄₁ belongs to the orthorhombic system



Figure 1 Subsolidus phase relations of the Y₂O₃-SrO-CuO system.

with a = 11.470 nm, b = 13.402 nm and c = 3.939 nm. The results obtained here are very close to those of McCarron *et al.* [13].

About 90 samples with various compositions of yttrium, strontium and copper were made. From X-ray diffraction measurements, the subsolidus phase relations in the Y_2O_3 -Sr-CuO system at 960 °C were determined. They are shown in Fig. 1, from which we find that in this system, there is no evidence for the existence of a ternary $YSr_2Cu_3O_7$ oxide. The section at room temperature in the Y_2O_3 -SrO-CuO system can be divided into the following sections:

(A) $Y_2O_3 + Y_2SrO_4 + Y_2Cu_2O_5$ (B) $Y_2SrO_4 + Sr_2CuO_3 + SrO$ (C) $Y_2SrO_4 + Sr_2CuO_3 + SrCuO_2$ (D) $Y_2SrO_4 + SrCuO_2 + Y_xSr_{14-x}Cu_{24}O_{41}$ (x = 4) (E) $Y_2SrO_4 + Y_2Cu_2O_5 + Y_xSr_{14-x}O_{41}(x = 4)$

(F) $SrCuO_2 + Y_xSr_{14-x}Cu_{24}O_{41}(0 < x < 4)$ (G) $CuO + Y_xSr_{14-x}Cu_{24}O_{41}(0 < x < 4)$

(H)
$$Y_2Cu_2O_5 + CuO + Y_rSr_{14-r}Cu_{24}O_{41}(x = 4)$$

For the compound $Sr_{14}Cu_{24}O_{41}$, Sr^{2+} can be par-tially replaced by Y^{3+} to form solid solution $Y_x Sr_{14-x} Cu_{24} O_{41}$ (0 < x \leq 4). The layered structure of Sr₁₄Cu₂₄O₄₁ shows that the structure includes both Cu-O planes and chains separated by strontium, and the copper ion in both structural elements is located at the centre of the squares of oxygen lying parallel to the a-c plane [13]. Furthermore, X-ray analysis shows that the solid solutions $Y_x Sr_{14-x} Cu_{24} O_{41}$ (0 < x \leq 4) and $Sr_{14} Cu_{24} O_{41}$ have isomorphous structure. In the solid solution $Y_x Sr_{14-x} Cu_{24} O_{41}$ (0 < x \leq 4), Y^{3+} occupies the same site as that of Sr^{2+} in $Sr_{14}Cu_{24}O_{41}$. The elements of the structure in $Y_x Sr_{14-x} Cu_{24}O_{41}(0 < x \le 4)$ are depicted in Fig. 2. Because the Y^{3+} ionic radius is less than that of the Sr^{2+} ion, replacing Sr^{2+} by Y^{3+} must change the lattice parameters. When the yttrium content increases, the variation in lattice parameters in $Y_x Sr_{14-x} Cu_{24} O_{41}$ (0 < x \leq 4) is as shown in Fig. 3. Because the structure is layered, the substitution dir-





Figure 2 The structure of $Y_xSr_{14-x}Cu_{24}O_{41}$ ($O < x \le 4$). The structure can be generated by stacking the layers (a-c) along the *b*-axis by simply aligning the unit cells outlined in the *a*-*c* plane. (o) Copper and (\bigcirc) oxygen ions. (a) The infinite [CuO_{4/2}] chains propagating along the *c*-axis. (b) The [Y_xSr_{14-x}] layer residing between the chains and sheets. (c) The [CuO_{3/3}O_{1/2}] sheets. Within the sheets O1 is the triple bridging oxygen; O2 is double bridging.

ectly affects the layer stacking axis (*b*-axis), while it has a lesser effect on the *a*-*c* plane parameters. The parameters of $Y_4Sr_{10}Cu_{24}O_{41}$ are a = 1.1360 nm, b = 1.2950 nm and c = 0.3938 nm.

Preliminary electrical resistivity measurements showed that $Sr_{14}Cu_{24}O_{41}$ is a semiconductor and $Y_xSr_{14-x}Cu_{24}O_{41}$ ($0 < x \le 4$) gradually becomes an insulator as the yttrium content increases. It is interesting to compare $Y_xSr_{14-x}Cu_{24}O_{41}$ ($0 < x \le 4$) with $YBa_2Cu_3O_{7-y}$. Structurally, they are layered and have as their basic squares planar copper atoms bound to oxygen. The square planar units are arranged so as to form both Cu–O chains and Cu–O planes. The difference between them appears to be in the mode of linkage of the square planar copper oxygen



Figure 3 Variation of the lattice parameters with yttrium content.

units: edge-shared in $Y_x Sr_{14-x} Cu_{24}O_{41}$ ($0 < x \le 4$), corner-shared in $YBa_2Cu_3O_{7-y}$. In addition, in $Y_x Sr_{14-x}Cu_{24}O_{41}$ ($0 < x \le 4$), there is no continuous corner-shared CuO₂ network, as there is in $YBa_2Cu_3O_{7-y}$, i.e. there is no true CuO₂ plane in this compound, therefore it does not exhibit superconductivity.

Acknowledgements

The authors thank Mr Ni Yongming and Mrs Jia Shunlian for electrical resistivity measurements.

References

- 1. K. A. MULLER and J. G. BEDNORZ, Z. Phys. B 64 (1986) 189.
- 2. C. W. CHU, et al., Phys. Rev. Lett. 58 (1987), 405.
- 3. ZHAO ZHONG-XIAN, CHEN LIQUAN, YANG QIAN-SHENG, et al., Kexue Tongbao 8 (1987) 522.
- 4. Idem, ibid. 10 (1987) 661.
- 5. Y. MEI, S. M. GREEN, C. JIAN and H. L. LUO, in Novel superconductivity", edited by S. A. Wolf and V. Z. Kresin (Plenum, New York, 1987), p. 1041.
- M. K. WU, J. R. ASHBURN, C. A. HIGGIN, et al., Phys. Rev. B. 37 (1988) 9765.
- M. ODA, T. MURAKAMI, ENOMOTO and M. SUZUKI, Jpn. J. Appl. Phys. 26 (1987) L804.
- ZHANG QI-RUI, CAO LIE-ZAO, QIAN YI-TAI, et al., Solid State Commun. 63 (1987) 535.
- 9. B. W. VEAL, W. K. KWOK, A. UMEZAWA et al., Appl. Phys. Lett. 51 (1987) 279.
- 10. S. G. TRESVYATSKII, L. M. LOPATO, A. E. KASHC-HEVSKII and A. V. SHEVCHENKO, *Inorg. Mater.* 7 (1971) 1614.
- 11. CHE GUANG-CAN, LIANG JING-KUI, CHEN WEI, XIE SI-SHEN, et al., Scientia Sinica A 16 (1988) 1079.
- 12. CHR. L. TESKE and H. MULLER-BUSCHBAUM, Z. Anorg. Allg. Chem. 371 (1969) 325.
- 13. Idem, ibid. 379 (1970) 324.
- E. M. McCARRON, M. A. SUBRAMANIAN, J. C. CALABRESE and R. L. HARLOW, Mater. Res. Bull. 23 (1988) 1355.

Received 12 February and accepted 20 June 1991